

Question 1:

a) Starting with **time harmonic** (sinusoidal steady state) form of the Maxwell's equations find the expressions for the electric field intensity and magnetic flux density in terms of the vector and scalar potentials,  $\vec{A}(r)$  and  $V(r)$ . Show all your work. [10 points]

b) Again, assuming time harmonic fields, derive the non-homogeneous wave equation for  $\vec{A}(r)$  and  $V(r)$  subject to the Lorentz gauge for a **simple medium**. Show all your work. [20 points]

a) From <sup>(1)</sup>  $\nabla \cdot \vec{B} = 0$  & <sup>(2)</sup>  $\text{div}(\text{curl of any vector}) = 0 \Rightarrow$

<sup>(3)</sup>  $\boxed{\vec{B} = \nabla \times \vec{A}}$  . From <sup>(4)</sup>  $\nabla \times \vec{E} = -j\omega \vec{B}$  & (3)  $\Rightarrow$

$\nabla \times \vec{E} = -j\omega(\nabla \times \vec{A}) \Rightarrow$  <sup>(5)</sup>  $\nabla \times \vec{E} + j\omega \nabla \times \vec{A} = 0 \Rightarrow$

<sup>(6)</sup>  $\nabla \times (\vec{E} + j\omega \vec{A}) = 0$  since  $\text{curl}(\text{grad of any scalar}) = 0 \Rightarrow$

$\vec{E} + j\omega \vec{A} = -\nabla V \Rightarrow$  <sup>(7)</sup>  $\boxed{\vec{E} = -\nabla V - j\omega \vec{A}}$

b) From Faraday's law <sup>(8)</sup>  $\nabla \times \vec{H} = \vec{J} + j\omega \epsilon \vec{E}$  or since  $\vec{H} = \frac{\vec{B}}{\mu}$  then

<sup>(9)</sup>  $\nabla \times \vec{B} = \mu \vec{J} + j\omega \mu \epsilon \vec{E}$  use (3) & (7) in (9)  $\Rightarrow$

$\nabla \times (\nabla \times \vec{A}) = \mu \vec{J} + j\omega \mu \epsilon (-\nabla V - j\omega \vec{A}) \Rightarrow$

$\nabla(\nabla \cdot \vec{A}) - \nabla^2 \vec{A} = \mu \vec{J} - j\omega \mu \epsilon \nabla V + \omega^2 \mu \epsilon \vec{A} \Rightarrow$

$-\nabla^2 \vec{A} = \mu \vec{J} - j\omega \mu \epsilon \nabla V - \nabla(\nabla \cdot \vec{A}) + \omega^2 \mu \epsilon \vec{A} \Rightarrow$

<sup>(10)</sup>  $-\nabla^2 \vec{A} = \mu \vec{J} - \nabla(j\omega \mu \epsilon V + \nabla \cdot \vec{A}) + \omega^2 \mu \epsilon \vec{A}$

(10)  $\Rightarrow$

(11)

$$-\nabla^2 \vec{A} = \mu \vec{J} + \omega^2 \mu \epsilon \vec{A} - \nabla (j\omega \mu \epsilon V + \nabla \cdot \vec{A})$$

For Lorenz's gauge  $j\omega \mu \epsilon V + \nabla \cdot \vec{A} = 0 \Rightarrow$

(12)

$$\nabla \cdot \vec{A} = -j\omega \mu \epsilon V \quad \text{then}$$

(11)  $\Rightarrow$  (13)  $-\nabla^2 \vec{A} = \mu \vec{J} + \omega^2 \mu \epsilon \vec{A}$  or

(14)

$$\nabla^2 \vec{A} + \omega^2 \mu \epsilon \vec{A} = -\mu \vec{J}$$

We now use the Gauss law to find the wave Eq for V

(15)

(16)  $\nabla \cdot \vec{D} = \rho \Rightarrow \nabla \cdot \vec{E} = \rho / \epsilon$  use (7) in (16)  $\Rightarrow$

$$\nabla \cdot (-\nabla V - j\omega \vec{A}) = \rho / \epsilon \Rightarrow \nabla^2 V - j\omega \nabla \cdot \vec{A} = \rho / \epsilon$$

but  $\nabla \cdot \vec{A}$  is given by (12)  $\Rightarrow$

$$-\nabla^2 V - j\omega (-j\omega \mu \epsilon V) = \rho / \epsilon \Rightarrow$$

$$-\nabla^2 V - \omega^2 \mu \epsilon V = \rho / \epsilon \Rightarrow$$

(18)

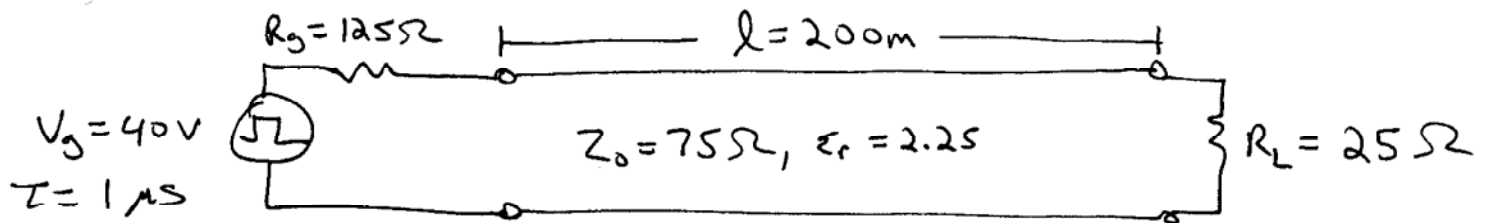
$$\nabla^2 V + \omega^2 \mu \epsilon V = -\rho / \epsilon$$

## Question 2

A  $75\ \Omega$  lossless transmission line, with  $\epsilon_r = 2.25$ , is connected to a  $40\text{ V}$  pulse generator with a source impedance of  $125\ \Omega$ . The line is  $200\text{ m}$  long and is terminated in a  $25\ \Omega$  load. At time  $t = 0$ , a single rectangular pulse of width  $1\ \mu\text{s}$  is sent down the line.

- Draw a reflection diagram for  $0 \leq t \leq 4\ \mu\text{s}$ . Be sure to include both the leading and trailing edges. Label each line with a numeric value. (15 marks)
- Plot  $V(z = 100\text{ m}, t)$  for  $0 \leq t \leq 4\ \mu\text{s}$ . Label all voltage levels and transition times with numeric values. (20 marks)

The TL circuit is



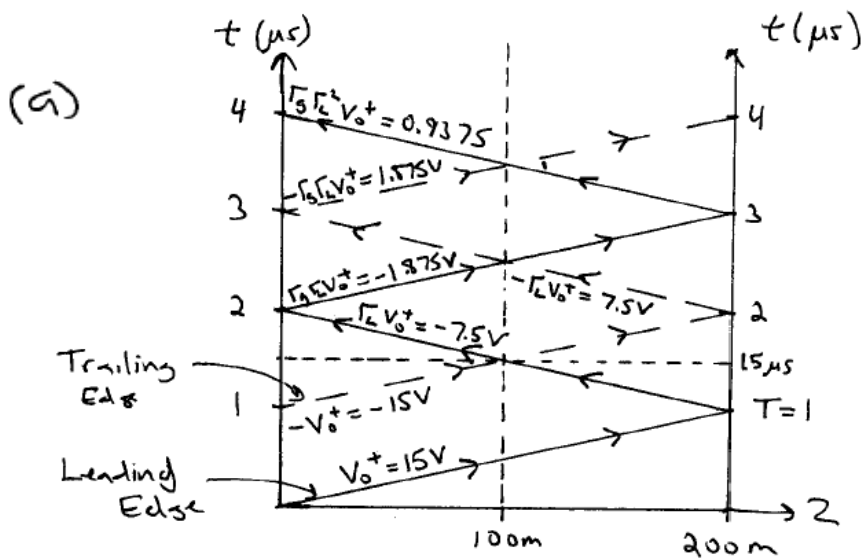
$$\text{For this line, } v_p = \frac{c}{\sqrt{\epsilon_r}} = \frac{3 \times 10^8}{\sqrt{2.25}} = 2 \times 10^8 \text{ m/s}$$

$$\therefore T = \frac{l}{v_p} = \frac{200}{2 \times 10^8} = \underline{\underline{1\ \mu\text{s}}}$$

$$V_o^+ = V_g \frac{Z_o}{Z_o + R_s} = 40 \frac{75}{75 + 125} = \underline{\underline{15\text{ V}}}$$

$$\Gamma_L = \frac{R_L - Z_o}{R_L + Z_o} = \frac{25 - 75}{25 + 75} = \underline{\underline{-\frac{1}{2}}}$$

$$\Gamma_s = \frac{R_s - Z_o}{R_s + Z_o} = \frac{125 - 75}{125 + 75} = \underline{\underline{\frac{1}{4}}}$$



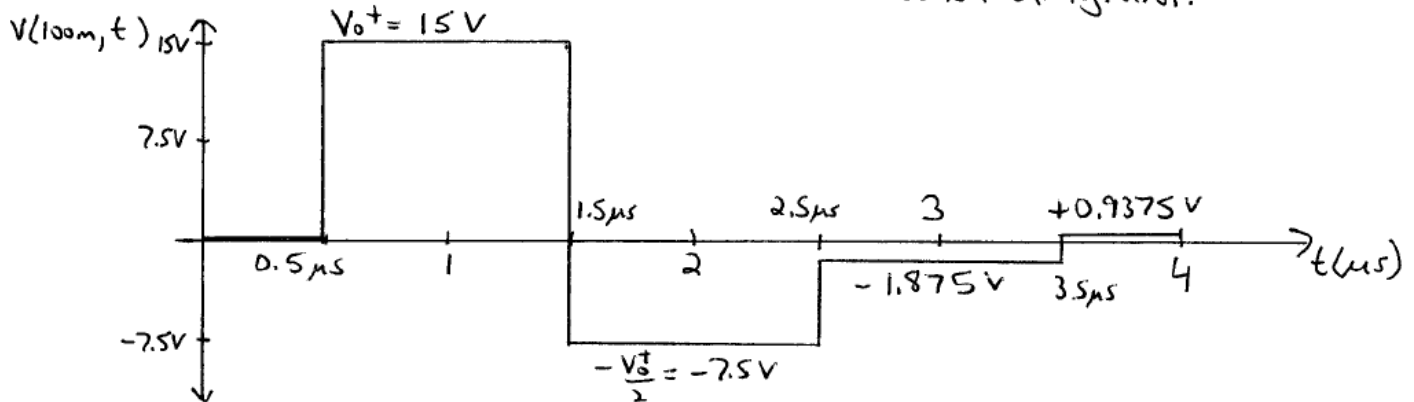
$$V_0^+ = 15\text{V}$$

$$\Gamma_L V_0^+ = \left(-\frac{1}{2}\right)(15) = -7.5\text{V}$$

$$\Gamma_S \Gamma_L V_0^+ = \left(\frac{1}{4}\right)(-7.5) = -1.875\text{V}$$

$$\Gamma_S \Gamma_L^2 V_0^+ = \left(-\frac{1}{2}\right)(-1.875) = 0.9375\text{V}$$

(b) To plot the Voltage versus time at  $z = 100\text{m}$ , follow dashed line at  $z = 100\text{m}$  in the reflection diagram:

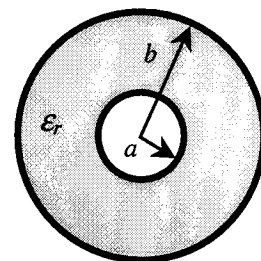




### Question 3

A coaxial cable lies along the  $z$ -axis, and has an inner conductor radius of  $\rho = a = 0.25$  mm, and an outer conductor radius of  $\rho = b = 0.75$  mm. The two conductors are separated by a non-magnetic ( $\mu_r = 1$ ), lossless dielectric, with  $\epsilon_r = 2$ . It is known, that the instantaneous electric field within this dielectric is given by

$$\vec{E} = \frac{200}{\rho} \sin(\omega t - \beta z) \hat{a}_\rho \text{ V/m for } a < \rho < b$$



with  $f = 30$  GHz ( $30 \times 10^9$  Hz).

- Find the instantaneous magnetic field,  $\vec{H}$ , within the cable, meaning for  $a < \rho < b$ . Show your work. (12 marks)
- Determine the expressions for the instantaneous surface charge density,  $\rho_s$ , and the instantaneous surface current density,  $\vec{J}_s$ , on the outer conductor ( $\rho = b$ ). You may assume that this is a perfect conductor. (10 marks)
- Describe the type of electromagnetic wave that these fields represent. Justify your answer. (4 marks)
- Find the numerical value of the propagation constant  $\beta$ . (4 marks)

a)  $\vec{H}$  can be found from Faraday's Law  $\nabla \times \vec{E} = -j\omega\mu_0\vec{H}$  ①  
but  $\vec{E}$  must first be written in phasor form:

$$\vec{E}(\rho, z, t) = \frac{200}{\rho} \sin(\omega t - \beta z) \hat{a}_\rho = \frac{200}{\rho} \cos(\omega t - \beta z - \frac{\pi}{2}) \hat{a}_\rho$$

$$\therefore \vec{E}(\rho, z) = \frac{200}{\rho} e^{-j\beta z} \underbrace{e^{-j\frac{\pi}{2}}}_{-j} \hat{a}_\rho = -j \frac{200}{\rho} e^{-j\beta z} \hat{a}_\rho$$

$$\therefore \text{From ①, } \vec{H} = \frac{j}{\omega\mu_0} \nabla \times \vec{E} = \left( \frac{j}{\omega\mu_0} \right) \left( \frac{\partial E_\rho}{\partial z} \hat{a}_\phi \right) \quad \left[ \begin{array}{l} \text{The curl has been} \\ \text{simplified in cylindrical} \\ \text{coordinates for } \vec{E}(\rho, z) \end{array} \right]$$

$$= \left( \frac{j}{\omega\mu_0} \right) \frac{\partial}{\partial z} \left[ -j \frac{200}{\rho} e^{-j\beta z} \right] \hat{a}_\phi$$

$$= \left( \frac{j}{\omega\mu_0} \right) \left( \frac{-j200}{\rho} \right) (-j\beta) \left( e^{-j\beta z} \right) \hat{a}_\phi$$

$$\vec{H}(\rho, z) = -j \frac{\beta}{\omega\mu_0} \left( \frac{200}{\rho} \right) e^{-j\beta z} \hat{a}_\phi = -j \frac{1}{\eta} \left( \frac{200}{\rho} \right) e^{-j\beta z} \hat{a}_\phi$$

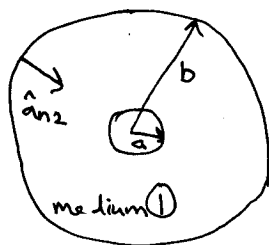
$\therefore$  The instantaneous  $\vec{H}$  field is then:

$$\hookrightarrow \text{Since } \eta = \frac{\omega\mu_0}{\beta} = \sqrt{\frac{\mu}{\epsilon}} = 266 \Omega$$

$$\vec{H}(\rho, z, t) = \text{Re}[\vec{H}(\rho, z) e^{j\omega t}] = \frac{1}{\eta} \frac{200}{\rho} \sin(\omega t - \beta z) \hat{a}_\phi = \underline{\underline{\frac{0.75}{\rho} \sin(\omega t - \beta z) \hat{a}_\phi \left( \frac{\text{A}}{\text{m}} \right)}}$$

(b) To solve for  $\rho_s$  &  $\vec{J}_s$  we must use boundary conditions:

medium ②



$$\Rightarrow \hat{a}_{n2} \cdot (\vec{D}_1 - \vec{D}_2) = \rho_s \quad \text{and} \quad \hat{a}_{n2} \times (\vec{H}_1 - \vec{H}_2) = \vec{J}_s$$

\* Our outward normal of medium 2 (Perfect conductor) (PEC) is  $\hat{a}_{n2} = -\hat{a}_\rho$ .

\* Also  $\vec{D}_2 = \vec{E}_2 = \vec{H}_2 = \vec{B}_2 = 0$  since it is a PEC.

$$\begin{aligned} \therefore \rho_s &= -\hat{a}_\rho \cdot \vec{D}_1(\rho=b, z, t) = -\hat{a}_\rho \cdot \epsilon_r \epsilon_0 \left[ \frac{200}{0.75 \text{ mm}} \sin(\omega t - \beta z) \hat{a}_\rho \right] \\ &= -\left( \frac{200}{0.75 \times 10^{-3}} \right) (2\epsilon_0) \sin(\omega t - \beta z) \\ \rho_s &= \underline{\underline{-4.7 \sin(\omega t - \beta z) \mu\text{C}/\text{m}^2}} \end{aligned}$$

$$\begin{aligned} \text{For } \vec{J}_s: \quad \vec{J}_s &= -\hat{a}_\rho \times \vec{H}_1(\rho=b, z, t) \\ &= -\hat{a}_\rho \times \left[ \frac{0.75}{0.75 \times 10^{-3}} \sin(\omega t - \beta z) \hat{a}_\phi \right] \\ \vec{J}_s &= \underline{\underline{-1000 \sin(\omega t - \beta z) \hat{a}_z \text{ A}/\text{m}^2}} \end{aligned}$$

(c) These fields,  $\vec{E}$  &  $\vec{H}$ , correspond to a TEM wave.

This is because both  $\vec{E}$  and  $\vec{H}$  are perpendicular to the direction of propagation.

(d) For a TEM wave,  $\beta = \omega \sqrt{\mu \epsilon}$

$$\begin{aligned} \therefore \beta &= 2\pi f \sqrt{\mu_r \epsilon_r \epsilon_0} \\ &= 2\pi(30 \times 10^9) \sqrt{\mu_0 (2) \epsilon_0} \\ \beta &= \underline{\underline{889 \text{ rad}/\text{m}}} \end{aligned}$$