

**HW #5**

Q1) A vector field  $\vec{A} = \hat{a}_P(3\cos\phi) - \hat{a}_\theta 2P + 5\hat{a}_z$  is expressed in cylindrical coordinates. Find its representation in rectangular coordinate system.

Sol: ①  $\begin{pmatrix} Ax \\ Ay \\ Az \end{pmatrix} = \begin{pmatrix} \cos\phi & -\sin\phi & 0 \\ \sin\phi & \cos\phi & 0 \\ 0 & 0 & 1 \end{pmatrix} \begin{pmatrix} Ap \\ A_\theta \\ Az \end{pmatrix} \Rightarrow$

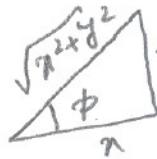
②  $Ax = \cos\phi Ap - \sin\phi A_\theta = \cos\phi(3\cos\phi) - \sin\phi(-2P) = 3\cos^2\phi + 2P\sin\phi$

③  $Ay = \sin\phi Ap + \cos\phi A_\theta = \sin\phi(3\cos\phi) + \cos\phi(-2P) = 3\cos\phi\sin\phi - 2P\cos\phi$

④  $Az = Az = 5$

\* Now we need to know how to convert the  $\cos\phi$ ,  $\sin\phi$  &  $P$  to rectangular coordinates. Recall

$$\begin{aligned} x &= P\cos\phi & x^2 + y^2 &= P^2 \\ y &= P\sin\phi & \tan\phi &= y/x \end{aligned}$$



$$\cos\phi = \frac{x}{\sqrt{x^2+y^2}}$$

$$\sin\phi = \frac{y}{\sqrt{x^2+y^2}}$$

$$P = \sqrt{x^2+y^2}$$

$$Ax = 3\cos^2\phi + 2P\sin\phi = 3 \frac{x^2}{x^2+y^2} + 2\sqrt{x^2+y^2} \frac{y}{\sqrt{x^2+y^2}} \Rightarrow$$

$$Ax = \boxed{\frac{3x^2}{x^2+y^2} + 2y}$$

$$Ay = 3\cos\phi\sin\phi - 2P\cos\phi = 3 \frac{x}{\sqrt{x^2+y^2}} \frac{y}{\sqrt{x^2+y^2}} - 2\sqrt{x^2+y^2} \frac{x}{\sqrt{x^2+y^2}} \Rightarrow$$

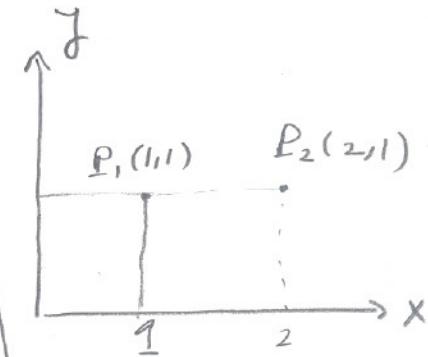
$$Ay = \boxed{\frac{3xy}{x^2+y^2} - 2x}$$

$$\text{ & } \boxed{Az = 5} \text{ then }$$

$$\vec{A} = Ax \hat{a}_x + Ay \hat{a}_y + Az \hat{a}_z = \left( \frac{3x^2}{x^2+y^2} + 2y \right) \hat{a}_x + \left( \frac{3xy}{x^2+y^2} - 2x \right) \hat{a}_y + 5 \hat{a}_z$$

Q2: Consider the central force  $\vec{F} = \frac{1}{r^2} \hat{a}_r$ . Calculate the work done in presence of the above field in moving from  $P_1(1,1)$  to  $P_2(2,1)$  along the line  $P_1P_2$

Sol: We convert the  $\vec{F} = \frac{1}{r^2} \hat{a}_r$  to rectangular



Coordinates:

$$\begin{pmatrix} F_x \\ F_y \\ F_z \end{pmatrix} = \begin{pmatrix} \sin\theta \cos\phi & \cos\theta \cos\phi & -\sin\phi \\ \sin\theta \sin\phi & \cos\theta \sin\phi & \cos\phi \\ \cos\theta & -\sin\theta & 0 \end{pmatrix} \begin{pmatrix} F_r \\ F_\theta = 0 \\ F_\phi = 0 \end{pmatrix} \Rightarrow$$

$$F_x = \sin\theta \cos\phi F_r$$

$$F_y = \sin\theta \sin\phi F_r$$

$$F_z = \cos\theta F_r$$

here since we are only in two-dimensional, i.e. we are in  $x-y$  plane for which  $\theta = \pi/2$   
we have

$$F_x = \cos\phi F_r = \cos\phi \frac{1}{r^2}$$

$$F_y = \sin\phi F_r = \sin\phi \frac{1}{r^2}$$

$$F_z = 0$$

\* The relation between  $(x, y, z)$  &  $(r, \theta, \phi)$  in 3-D is given by

$$x = r \sin\theta \cos\phi$$

$$y = r \sin\theta \sin\phi$$

$$z = r \cos\theta$$

In 2D (-x-y Plane)  $\theta = \pi/2 \Rightarrow$

$$x = r \cos\phi$$

$$y = r \sin\phi$$

$$z = 0$$

$$x^2 + y^2 = r^2$$

$$\cos\phi = \frac{x}{r} = \frac{x}{\sqrt{x^2 + y^2}}$$

$$\sin\phi = \frac{y}{r} = \frac{y}{\sqrt{x^2 + y^2}}$$

then

$$F_x = c_0 + \frac{1}{r^2} = \frac{x}{\sqrt{x^2+y^2}} \quad \frac{1}{x^2+y^2} = \frac{x}{(x^2+y^2)^{3/2}}$$

$$F_y = \sin \phi \frac{1}{r^2} = \frac{y}{\sqrt{x^2+y^2}} \quad \frac{1}{x^2+y^2} = \frac{y}{(x^2+y^2)^{3/2}}$$

$$\vec{F} = F_x \hat{ax} + F_y \hat{ay} = \frac{1}{(x^2+y^2)} \left[ \frac{x}{\sqrt{x^2+y^2}} \hat{ax} + \frac{y}{\sqrt{x^2+y^2}} \hat{ay} \right]$$

$$W = \int \vec{F} \cdot d\vec{l} = \int (F_x \hat{ax} + F_y \hat{ay}) \cdot (dx \hat{ax} + dy \hat{ay})$$

$$= \int F_x dx + F_y dy$$

but in going from  $P_1$  to  $P_2$  we move along the line  $y=1 \Rightarrow dy=0$

$$W = \int F_x dx = \int_{x=1}^{x=2} \frac{x}{(x^2+y^2)^{3/2}} dx$$

We need to express  $y$  in terms of  $x$ . Note that in going from  $P_1$  to  $P_2$   $y$  is constant & equal to 1  $\Rightarrow$

$$W = \int_{x=1}^{x=2} \frac{x}{(x^2+1)^{3/2}} dx$$

$$\text{let } U = x^2+1 \quad \frac{1}{2} du = x dx \Rightarrow \frac{1}{2} du = x dx \Rightarrow$$

$$du = 2x dx \Rightarrow$$

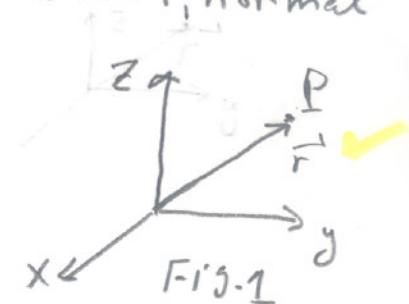
$$W = \int \frac{1}{2} \frac{du}{U^{3/2}} = \frac{1}{2} \left( \frac{U^{-1/2}}{-1/2} \right) =$$

$$-\left( \frac{1}{\sqrt{U}} \right) \Big|_{x=1}^{x=2} = \frac{1}{\sqrt{2}} - \frac{1}{\sqrt{5}}$$

Q3 ✓

a) Let  $\vec{r}$  be the position vector, pointing from origin to the point of observation  $P$ . show that

$\nabla |\vec{r}| = \nabla r = \frac{\vec{r}}{|\vec{r}|} = \hat{r}$  where  $\hat{r}$  is the unit normal in the direction of  $\vec{r}$ .



b) let function  $g(r)$  depend only on the distance

between the source points (Primed coordinates)

& the observation point (Unprimed coordinates), see Fig 2.

prove that  $\nabla g(r) = Df(R) = \hat{r}_R \frac{\partial f(R)}{\partial R}$ . use this information to calculate the gradient of  $g(x, y, z) = f(R) = \sin\left(\pi \frac{R^2}{4}\right)$

Sol:

$$\text{a) } \vec{r} = x \hat{a}_x + y \hat{a}_y + z \hat{a}_z$$

$$|\vec{r}| = r = [x^2 + y^2 + z^2]^{1/2}$$

$$\begin{aligned} \nabla |\vec{r}| = \nabla r &= \frac{2x}{2\sqrt{x^2+y^2+z^2}} \hat{a}_x + \\ &\quad \frac{2y}{2\sqrt{x^2+y^2+z^2}} \hat{a}_y + \frac{2z}{2\sqrt{x^2+y^2+z^2}} \hat{a}_z \\ &= \frac{x \hat{a}_x + y \hat{a}_y + z \hat{a}_z}{\sqrt{x^2+y^2+z^2}} = \frac{\vec{r}}{|\vec{r}|} = \hat{r} \end{aligned}$$

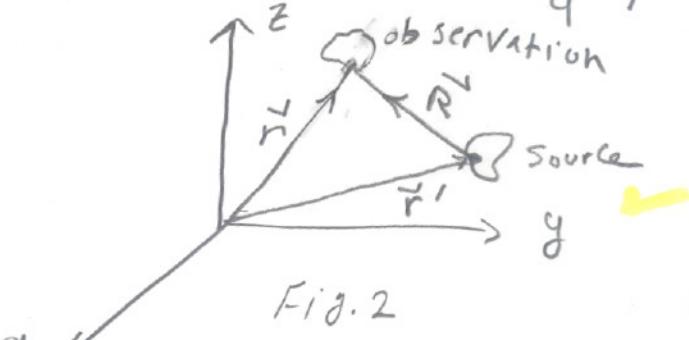


Fig. 2

b) we note the following ~~we note that~~

5

$$\vec{R} = (x-x') \hat{a}_x + (y-y') \hat{a}_y + (z-z') \hat{a}_z \quad \&$$

$$R = |\vec{R}| = [(x-x')^2 + (y-y')^2 + (z-z')^2]^{1/2} \text{ then}$$

$$\nabla |\vec{R}| = \nabla R = \frac{(x-x') \hat{a}_x + (y-y') \hat{a}_y + (z-z') \hat{a}_z}{\sqrt{(x-x')^2 + (y-y')^2 + (z-z')^2}}$$
$$= \frac{\vec{R}}{|\vec{R}|} = \hat{a}_R$$

Also from chain rule for differentiation

$$\nabla g(r) = \nabla f(R) = \frac{\partial f(R)}{\partial R} \nabla R = \frac{\partial f(R)}{\partial R} \hat{a}_R$$

using the above to calculate

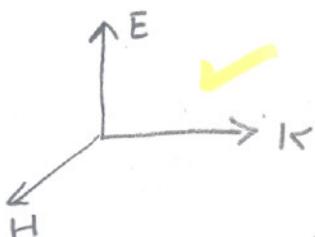
$$\nabla f(R) = \nabla \left[ \sin \left( \pi \frac{R^2}{4} \right) \right] = \frac{\partial}{\partial R} \sin \left( \pi \frac{R^2}{4} \right) \hat{a}_R$$
$$= \frac{2\pi R}{4} \cos \left( \pi \frac{R^2}{4} \right) \hat{a}_R = \boxed{\frac{\pi R}{2} \cos \left( \pi \frac{R^2}{4} \right) \hat{a}_R}$$

Q4: Consider a plane wave  $\vec{E} = \vec{E}_0 e^{-j\vec{k} \cdot \vec{r} + j\omega t}$  propagating in a homogeneous, lossless, source free region for which  $\epsilon > 0$  &  $\mu > 0$  &  $\vec{E}_0$  is constant. b

a) Show that  $\vec{k} \perp \vec{E}$  &  $\vec{k} \perp \vec{H}$

b) Show that  $\vec{k}$ ,  $\vec{E}$  &  $\vec{H}$  form a right hand triplet (Hint

$$\vec{k} \times \vec{E} = \omega \mu \vec{H} \quad \& \quad \vec{k} \times \vec{H} = -\epsilon \omega \vec{E}$$



Sol: since source free <sup>①</sup>  $\nabla \cdot \vec{E} = 0$  &

$$\nabla \cdot \vec{B} = 0$$

③

$$\nabla \cdot (\vec{F} \vec{F}) = (\nabla \cdot \vec{F}) \vec{F} + \vec{F} \cdot \nabla \vec{F}$$

From vector calculus  $\nabla \cdot (\vec{F} \vec{F}) = (\nabla \cdot \vec{F}) \vec{F} + \vec{F} \cdot \nabla \vec{F}$   
then <sup>④</sup>  $\nabla \cdot (\vec{E}_0 e^{-j\vec{k} \cdot \vec{r} + j\omega t}) = (\nabla \cdot \vec{E}_0) e^{-j\vec{k} \cdot \vec{r} + j\omega t} + \nabla (e^{-j\vec{k} \cdot \vec{r} + j\omega t}) \cdot \vec{E}_0$

since  $\vec{E}_0$  is constant <sup>⑤</sup>  $\nabla \cdot \vec{E}_0 = 0 \Rightarrow$

$$\nabla \cdot (\vec{E}_0 e^{-j\vec{k} \cdot \vec{r} + j\omega t}) = \nabla (e^{-j\vec{k} \cdot \vec{r} + j\omega t}) \cdot \vec{E}_0 = -j\vec{k} e^{-j\vec{k} \cdot \vec{r} + j\omega t} \cdot \vec{E}_0 = -j(\vec{k} \cdot \vec{E}_0) e^{-j\vec{k} \cdot \vec{r} + j\omega t}$$

but from (1) we must have

$$\nabla \cdot (\vec{E}_0 e^{-j\vec{k} \cdot \vec{r} + j\omega t}) = 0 \Rightarrow \vec{k} \cdot \vec{E}_0 = 0 \Rightarrow \boxed{\vec{k} \perp \vec{E}_0 \Rightarrow \vec{k} \perp \vec{E}}$$
⑧

\* Using equation (2) & similar argument, i.e

$$\nabla \cdot \vec{B} = 0 \Rightarrow \nabla \cdot \vec{H} = 0 \Rightarrow \nabla \cdot (\vec{H}_0 e^{-j\vec{k} \cdot \vec{r} + j\omega t}) = 0 \Rightarrow \vec{k} \cdot \vec{H}_0 = 0 \Rightarrow \boxed{\vec{k} \perp \vec{H}_0 \Rightarrow \vec{k} \perp \vec{H}}$$
⑩

b) since source free & simple medium

(12)

$$(11) \nabla \times \vec{E} = -\frac{\partial \vec{B}}{\partial t} \xrightarrow{\text{Time-Harmonics}} \boxed{\nabla \times \vec{E} = -j\omega \vec{B} = -j\omega \mu \vec{H}} \quad \text{Z}$$

$$(13) \nabla \times \vec{H} = +\frac{\partial \vec{D}}{\partial t} \xrightarrow{\text{Time-Harmonics}} \boxed{\nabla \times \vec{H} = j\omega \vec{D} = j\omega \epsilon \vec{E}} \quad (14)$$

Let's calculate (15)  $\nabla \times \vec{E} = \nabla \times (\vec{E}_0 \vec{e}^{j\vec{k} \cdot \vec{r} + j\omega t})$

The following identity holds

$$(16) \nabla \times (f \vec{F}) = \nabla f \times \vec{F} + f \nabla \times \vec{F}$$

Using (16) in (15)

$$\begin{aligned} \nabla \times \vec{E} &= \nabla \times (\vec{E}_0 \vec{e}^{j\vec{k} \cdot \vec{r} + j\omega t}) = \\ &= (\nabla \vec{e}^{j\vec{k} \cdot \vec{r} + j\omega t}) \times \vec{E}_0 + \vec{e}^{j\vec{k} \cdot \vec{r} + j\omega t} \cancel{\nabla \times \vec{E}_0} \\ &= -j\vec{k} \vec{e}^{j\vec{k} \cdot \vec{r} + j\omega t} \times \vec{E}_0 = -j\vec{k} \times \vec{E}_0 \vec{e}^{-j\vec{k} \cdot \vec{r} + j\omega t} \Rightarrow \end{aligned}$$

$$\boxed{\nabla \times \vec{E} = -j\vec{k} \times \vec{E}} \quad (17)$$

Sub (17) in (12)  $\Rightarrow \cancel{-j\vec{k} \times \vec{E}} = -j\omega \mu \vec{H} \Rightarrow$

$$(18) \boxed{\vec{k} \times \vec{E} = \omega \mu \vec{H}}$$

\* We now use  $\vec{H} = \vec{H}_0 \vec{e}^{-j\vec{k} \cdot \vec{r} + j\omega t}$  to calculate  $\nabla \times \vec{H}$  & use the result in (14) to obtain the following

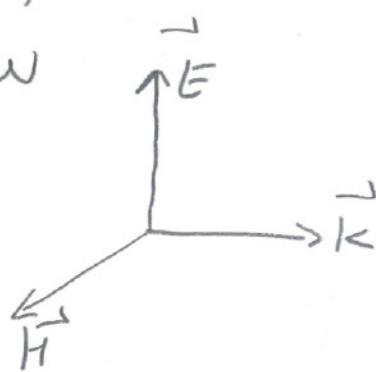
$$\textcircled{19} \quad -j \vec{K} \times \vec{H} = j\omega \epsilon \vec{E} \Rightarrow \textcircled{20} \quad \boxed{\vec{K} \times \vec{H} = -\omega \epsilon \vec{E}}$$

8

To summarize we have shown

$$\vec{K} \perp \vec{E}, \vec{K} \perp \vec{H} \text{ & } \vec{K} \times \vec{E} = \omega \mu \vec{H} \text{ & } \vec{K} \times \vec{H} = -\omega \epsilon \vec{E}$$

from these we see that  $\vec{K}, \vec{E} \text{ & } \vec{H}$  Form a Right Hand triplet as shown below



Q5: Starting with differential form of the Maxwell's equations, find their integral form formulation.  
sol:

\*  $\nabla \times \vec{E} = -\frac{\partial \vec{B}}{\partial t}$  Faraday's law

$$\iint_S \nabla \times \vec{E} \cdot d\vec{s} = -\frac{\partial}{\partial t} \iint_S \vec{B} \cdot d\vec{s} \text{ use Stokes theorem}$$

$$\boxed{\oint_C \vec{E} \cdot d\vec{l} = -\frac{\partial}{\partial t} \Phi_B} \text{ where } \Phi_B = \iint_S \vec{B} \cdot d\vec{s}$$

\*  $\nabla \times \vec{H} = \frac{\partial \vec{D}}{\partial t} + \vec{J}$  Ampere's law. Use Stokes theorem

$$\iint_S \nabla \times \vec{H} \cdot d\vec{s} = \frac{\partial}{\partial t} \iint_S \vec{D} \cdot d\vec{s} + \iint_S \vec{J} \cdot d\vec{s} \Rightarrow$$

$$\boxed{\oint_C \vec{H} \cdot d\vec{l} = \frac{\partial}{\partial t} \iint_S \vec{D} \cdot d\vec{s} + I} \text{ where } I = \iint_S \vec{J} \cdot d\vec{s}$$

\*  $\nabla \cdot \vec{D} = \rho_e$  Gauss law. Use divergence theorem

$$\iiint_V \nabla \cdot \vec{D} dv = \iiint_V \rho_e dv \Rightarrow$$

$$\boxed{\iint_S \vec{D} \cdot d\vec{s} = Q_{enc}} \text{ where } Q_{enc} = \iiint_V \rho_e dv$$

\*  $\nabla \cdot \vec{B} = 0$  use divergence theorem

$$\iiint_V \nabla \cdot \vec{B} dv = 0 \Rightarrow \boxed{\iint_S \vec{B} \cdot d\vec{s} = 0}$$